

## Early Journal Content on JSTOR, Free to Anyone in the World

This article is one of nearly 500,000 scholarly works digitized and made freely available to everyone in the world by JSTOR.

Known as the Early Journal Content, this set of works include research articles, news, letters, and other writings published in more than 200 of the oldest leading academic journals. The works date from the mid-seventeenth to the early twentieth centuries.

We encourage people to read and share the Early Journal Content openly and to tell others that this resource exists. People may post this content online or redistribute in any way for non-commercial purposes.

Read more about Early Journal Content at <a href="http://about.jstor.org/participate-jstor/individuals/early-journal-content">http://about.jstor.org/participate-jstor/individuals/early-journal-content</a>.

JSTOR is a digital library of academic journals, books, and primary source objects. JSTOR helps people discover, use, and build upon a wide range of content through a powerful research and teaching platform, and preserves this content for future generations. JSTOR is part of ITHAKA, a not-for-profit organization that also includes Ithaka S+R and Portico. For more information about JSTOR, please contact support@jstor.org.

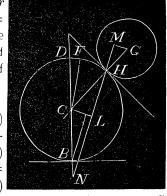
## SOLUTION OF PROBLEM 89. (SEE PAGE 195, VOL. II.)

## BY R. J. ADCOCK, MONMOUTH, ILLINOIS.

LET m, m', be the masses of the spheres R and r; x, y, the horizontal and vertical coordinates of r's centre with reference to R's centre when in its initial position,  $\theta = DCG =$  angle between line of centres and vertical at

time t,  $\theta_1$ ,  $\theta_2$  = angular rotations of R and r about centres,  $\alpha$  = initial value of  $\theta$ , F = magnitude and direction of the resultant of the reaction between the spheres at the point of contact H, p = the magnitude and direction of the normal pressure at H.

From the geometrical relations,  $x = R\theta_1 + (R+r)\sin\theta\dots(1)$ ,  $y = (R+r)\cos\theta,\dots(2)$   $R(\theta-\alpha-\theta_1) = \text{extent of arcs bro't into contact}$ ,  $R(\theta-\alpha-\theta_1) \div r + \theta - \alpha = [(R+r)(\theta-\alpha) - R\theta_1] \div r = \theta_2\dots(3)$ ,  $dx^2 + dy^2 = (R+r)d\theta^2 + R^2d\theta_1^2 + 2R(R+r)\cos\theta d\theta d\theta_1$ , (4)



There being no loss of vis viva in perfect rolling,  $2m'g(R+r)(\cos \alpha - \cos \theta)$   $= m' \frac{dx^2 + dy^2}{dt^2} + \frac{2}{5}m'r^2\frac{d\theta_2^2}{dt^2} + \frac{7}{5}mR^2\frac{d\theta_1^2}{dt^2} = \frac{7}{5}m'(R+r)^2\frac{d\theta_2^2}{dt^2} + \frac{7}{5}(m+m')R^2\frac{d\theta_1^2}{dt^2} + \frac{7}{5}(m+m')R^2\frac{d\theta_1^2}{dt^2} + 2R(R+r)(\cos \theta - \frac{2}{5})m'\frac{d\theta d\theta_1}{dt^2} \dots (5); F \sin \frac{p}{F} = \text{tangential component of}$   $F \text{ at } H, F \cos \frac{p}{F} = p = m'g \cos \theta - m'(R+r)\frac{d\theta^2}{dt^2} \dots (6), = \text{normal pressure at } H, = \text{normal component of weight of } r \text{ minus its centrifugal force.}$ 

 $F\sin\left(\theta - \frac{p}{F}\right)$  = horizontal component of F at H, acting to right on lower sphere, also = horizontal force at B, the lower point of contact. F, acting in direction HN, is the only force communicated by the weight of the upper sphere to the lower, and therefore the lower sphere will tend to move to the right or to the left according as HN passes to the right or left of B.

By the principle that the angular acceleration of a body about a fixed axis = the moment of the impressed forces divided by the moment of inertia with respect to that axis (Bartlett's An. Mech. 8th ed. p. 248), I have the following four eq'ns: The angular acceleration of r about its central axis = that which would be produced by the force F acting upward with the lever arm GM, while the body were retained by a fixed axis through G, hence

Fr 
$$\sin \frac{p}{F} \div \frac{2}{5}m'r^2 = \frac{d^2\theta_2}{dt^2}$$
...(7); FR  $\left(\sin \frac{p}{F} - \sin(\theta - \frac{p}{F})\right) \div \frac{7}{5}mR^2 = \frac{d^2\theta_1}{dt^2}$ ,.(8) for lower sphere about axis through B.

For common rotation of both spheres about axis through centre of lower,

$$FR\sin\left(\theta-\frac{p}{F}\right) \div \left\{ \frac{2}{5}mR^2 + m'\left[\frac{2}{5}r^2 + (R+r)^2\right] \right\} = \frac{d^2\theta_1}{dt^2}, \dots (9)$$

and for rotation of upper sphere about axis through H,

$$r\left(m'g\sin\theta - m'(R+r)\frac{d^{2}\theta_{1}}{dt^{2}}\right) \div \frac{7}{5}m'r^{2} = \frac{d^{2}\theta_{2}}{dt^{2}}..........(10)$$

Eliminating  $F \sin \frac{p}{F}$  and  $F \sin \left(\theta - \frac{p}{F}\right)$  from (8) by (7) and (9),

$$\frac{2}{5}m'r\frac{d^2\theta_2}{dt^2} = \frac{9mR^2 + m'[2r^2 + 5(R+r)^2]}{5R} \cdot \frac{d^2\theta_1}{dt^2} \cdot \dots$$
 (11)

Integrating twice and observing that  $\theta_1$  and  $\theta_2$  begin together,

$$\theta_2 = \frac{(9m + 5m')R^2 + 7m'r^2 + 10m'rR}{2m'rR} \times \theta_1....(12)$$

Hence by (3) 
$$\theta_1 = \frac{2m'R(r+R)(\theta-a)}{(9m+7m')R^2+7m'r^2+10m'rR}$$
....(13)

Substituting in (1)  $\theta_1$  from (13) and  $\theta$  from (2)

$$\begin{aligned} x &= 2m'R^2(R+r)(\theta-\alpha) \div \left[9mR^2 + m'(7R^2 + 7r^2 + 10rR)\right] + (R+r)\sin\theta \\ &= 2m'R^2(R+r) \left\{ \left[\cos^{-1}\frac{y \div (R+r)}{y \div (R+r)} - \alpha\right] \right\} \div \left[9mR^2 + m'(7R^2 + 7r^2 + 10rR)\right] + \sqrt{[(R+r)^2 - y^2] \dots (14)}, \text{ which is the required equation.} \end{aligned}$$

From the figure 
$$NC: R :: \sin CHN : \sin N$$
; hence by (7) and (9) 
$$NC: R :: \frac{2}{5}m'r\frac{d^2\theta_2}{dt^2} : \frac{\frac{2}{5}mR^2 + m'[\frac{2}{5}r^2 + (R+r)^2]}{R} \frac{d^2\theta_1}{dt^2}, \quad \text{and} \quad$$

by (11)  $NC: R:: 9mR^2 + m'\lceil 2r^2 + 5(R+r)^2 \rceil : 2mR^2 + m'\lceil 2r^2 + 5(R+r)^2 \rceil$ , in which the first term of second couplet being greater than the second term by  $7mR^2$ , NC is in all cases greater than R, and therefore F or HN passes to the right of the point of contact B, and both spheres in all cases have their motions on the same side of the origin\*.

The value of  $d\theta_1$  from (13) in (5) gives  $\frac{d\theta^2}{dt^2}$ , which in (6),  $p = m'g \cos \theta$ 

 $-(R+r)m'\frac{d\theta^2}{dt^2}=0$ , gives an equation of the 2nd degree in  $\cos\theta$ , from which one of the values of  $\cos \theta$  gives the point of separation.

By equality of (7) and (10), 
$$F \sin \frac{p}{F} = \frac{2}{7}m'g\sin\theta - \frac{2}{7}m'\frac{R+r}{r} \cdot \frac{d^2\theta_1}{dt^2}$$
,

which shows that the tangential component employed in giving rotation to upper sphere about its central axis  $=\frac{2}{7}$  of the tangential component of the weight m'g, only when the surface on which the rolling takes place has no motion.

<sup>\*</sup>We dissent from this conclusion, 1st: Because it may easily be shown, from a different course of reasoning, that, in certain positions, the spheres will roll in opposite directions;